An Upper Bound for the First Zero of Bessel Functions

By Ll. G. Chambers

Abstract. It is shown, using the Rayleigh-Ritz method of the calculus of variations, that an upper bound for the first zero j_{ν} , of $z^{-\nu} J_{\mu}(z)$, $\nu > -1$, is given by

$$(\nu + 1)^{1/2} \{ (\nu + 2)^{1/2} + 1 \},\$$

and that for large ν , $j_{\nu} = \nu + O(\nu^{1/2})$.

1. The following upper bound is given by Watson [4] for the first zero j_{ν} of $J_{\nu}(x)$ $(\nu > 0)$

(1)
$$j_{\nu} < \left\{\frac{4}{3}(\nu+1)(\nu+5)\right\}^{1/2}$$
.

It may be shown that a better bound may be obtained, valid for $\nu > -1$, namely

(2)
$$(\nu+1)^{1/2}((\nu+2)^{1/2}+1).$$

2. Consider the function

(3)
$$u(z) = \Gamma(\nu+1)(2/(\gamma z))^{\nu}J_{\nu}(\gamma z).$$

The differential equation satisfied by u(z) is given by Watson [3] to be

(4)
$$z^2u'' + (2\nu + 1)zu' + \gamma^2 z^2 u = 0$$

with the boundary condition u(0) = 1, and if γ is a zero of J_{ν} , u(1) = 0.

Equation (4) can be written in Sturm-Liouville form

(5)
$$\frac{d}{dz}\left(z^{2\nu+1}\frac{du}{dz}\right)+\gamma^2 z^{2\nu+1}u=0.$$

Multiplying Eq. (5) by u and integrating over $0 \le z \le 1$, it follows that

(6)
$$\gamma^{2} = \frac{\int_{0}^{1} z^{2\nu+1} u^{\prime 2} dz}{\int_{0}^{1} z^{2\nu+1} u^{2} dx}$$

On integration by parts, $uu'z^{2\nu+1}$ will vanish at z = 0, if $\nu > -\frac{1}{2}$, and u(1) vanishes. Thus the relation (6) provides a variational formulation, as indicated by Irving and Mullineux [1], for γ^2 which is an eigenvalue for the differential equation (5). The first eigenvalue will be j_{ν}^2 . The functional

(7)
$$\Lambda(\omega) = \frac{\int_0^1 z^{2\nu+1} \omega'^2 dz}{\int_0^1 z^{2\nu+1} \omega^2 dz},$$

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as indicated by Irving and Mullineux [2], obeys the following relations

(8a)
$$\Lambda(u) = j_{\nu}^2,$$

(8b) $\Lambda(\omega) > \Lambda(u), \quad u \neq \omega.$

Thus $\Lambda(\omega)$ provides an upper bound to j_{ν}^2 when $\omega(1) = 0$, by the Rayleigh-Ritz procedure.

3. Consider the approximating function

$$\omega = 1 - z^{\mu}$$

where p is as yet unspecified.

(10)
$$\Lambda(\omega) = \frac{\int_0^1 z^{2\nu+1} p^2 z^{2p-2} dz}{\int_0^1 (1 - 2z^p + z^{2p}) z^{2\nu+1} dz}$$

(11)
$$= p^{2} \left\{ \frac{1/(2\nu+2p)}{1/(2\nu+2) - 2/(2\nu+p+2) + 1/(2\nu+2p+2)} \right\}$$

(12)
$$= \frac{(\nu+1)(2\nu+p+2)(\nu+p+1)}{\nu+p}$$

on simplification.

Up till now p has not been specified. $\Lambda(\omega)$ may be regarded as a function of p, and the best upper bound will follow by minimizing $\Lambda(\omega)$ with respect to p. It can be verified by differentiation that this happens when

(13)
$$p + \nu = (\nu + 2)^{1/2}$$
.

Although p is negative outside of $-1 < \nu < 2$, the process is still valid because all of the denominators in the expression (11) remain positive, and it can easily be seen that

$$\lim_{z\to 0} z^{2\nu+1}\omega\omega'=0,$$

so that the endpoint condition at z = 0 remains satisfied. It follows that

$$j_{\nu}^{2} < \frac{(\nu+1)\left[(\nu+2)+(\nu+2)^{1/2}\right]\left[(\nu+2)^{1/2}+1\right]}{(\nu+2)^{1/2}},$$

which gives

(14)
$$j_{\nu} < (\nu + 1)^{1/2} ((\nu + 2)^{1/2} + 1).$$

A straightforward reduction shows that, if $\nu + 1 > 0$,

$$(\nu+1)^{1/2}((\nu+2)^{1/2}+1) < \left\{\frac{4}{3}(\nu+1)(\nu+5)\right\}^{1/2}$$

(For $\nu = 7$ there is in fact equality.)

It thus follows that [4]

(15)
$$\{\nu(\nu+2)\}^{1/2} < j_{\nu} < (\nu+1)^{1/2} ((\nu+2)^{1/2}+1), \quad \nu > 1.$$

It follows from (15) that

(16)
$$j_{\nu} = \nu + O(\nu^{1/2})$$
 for large ν .

As an example, the bound for $\nu = 0$ is given by $\sqrt{2} + 1 = 2.4142$ in comparison with the true value 2.4048.

(9)

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1. J. IRVING & N. MULLINEUX, Mathematics in Physics and Engineering, Academic Press, New York, 1959, p. 388.

2. J. IRVING & N. MULLINEUX, Loc. cit., p. 39.

3. G. N. WATSON, A Treatise on the Theory of Bessel Functions, The University Press, Cambridge, 1944, p. 98.

4. G. N. WATSON, Loc. cit., p. 486.